

CHAOTIC ORBITS OF TUMBLING ELLIPSOIDS IN VISCOUS AND INVISCID FLUIDS

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Kozlov & Oniscenko (1982) showed that general ellipsoids can show chaotic motion in inviscid flow

Objectives:

- Can chaotic behaviour be seen in viscous flows using DNS
- Can chaotic dynamics be used to enhance mixing.





Problem Set-up



Chaotic motion <u>only</u> under conditions of nonintegrability of Kirchhoff's equations (Kozlov & Oniscenko, Sov. Math. Dokl 1982)



General Ellipsoid: $a \neq b \neq c$ Domain Scale: L = 512*a* $Re_p = \frac{2Va}{v}$ E = $\frac{k_t}{k_r} = \frac{mV^2}{I\Omega^2}$



MOTION OF ELLIPSOID IN INVISCID FLOW

$$\begin{pmatrix} \boldsymbol{L} \\ \boldsymbol{P} \end{pmatrix} = \begin{pmatrix} I & D \\ D^T & M \end{pmatrix} \triangleq \mathcal{M} \begin{pmatrix} \boldsymbol{\omega} \\ \mathbf{v} \end{pmatrix}$$

$$\dot{L} = L \times \omega + P \times v + T_s$$
$$\dot{P} = P \times \omega + F_s$$

$$\boldsymbol{q} = (\boldsymbol{L}, \boldsymbol{P})$$
$$\mathcal{H}(\boldsymbol{q}) = \frac{1}{2} \boldsymbol{q} \cdot \mathcal{M}^{-1} \boldsymbol{q}$$

L – Angular Momentum

 ω – Angular Velocity

P – Linear Momentum v – Linear Velocity



The Kirchhoff equations can be view as a Lie algebra, SE(3)

Casimir functions are:

$$C_1 = \mathbf{P} \cdot L;$$
 $C_2 = \|P\|^2$

In the general case $C_2 \neq 0$

 \mathcal{H} restricts system to a level set (coadjoint orbit)

For the equations to be integrable additional conversed quantity is required



ELLIPSOID ORBITS IN INVISCID FLUID

Periodic



Chaos





RECURRENCE QUANTIFICATION ANALYSIS (RQA) OF ELLIPSOID MOTION



$$DET = \frac{\sum_{\ell=\ell_{min}}^{N} \ell P(\ell)}{\sum_{\ell=1}^{N} \ell P(\ell)}$$

$$R(i,j) = \begin{cases} 1 \ if \ L_{\infty}(\overrightarrow{x_{i}}, \overrightarrow{x_{j}}) \leq \varepsilon \\ 0 \ otherwise \end{cases}$$

The Lyapunov exponent and Rényi Entropy of the system encoded in recurrence plot.

Recurrence Rate (RR)

$$RR = \frac{1}{N^2} \sum_{i,j=1}^{N} R(i,j)$$

 $ENTR = -\sum_{\ell=\ell_{min}}^{N} p(\ell) \ln p(\ell)$

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ELLIPSOID ORBITS AT DIFFERENT ENERGY RATIOS

Periodic Orbit $E = 1 \frac{\rho_l}{\rho_s} = 0.125$ Chaotic Orbit $E = 20 \frac{\rho_l}{\rho_s} = 8$







TIME DEPENDENT RQA ANALYSIS

Periodic



Chaos





UNDERSTANDING CHAOS

Energy Transfer

Poincare Plots







PARAMETRIC MAPS







EFFECTS OF VISCOSITY ON ORBITS OF SYMMETRIC ELLIPSOIDS



Chaotic motion <u>only</u> under conditions of non-integrability of Kirchhoff's equations (Kozlov & Oniscenko, Sov. Math. Dokl 1982). Vortex shedding breaks the symmetry of the system, and allows chaotic orbits





VISCOSITY GRADIENTS AND MIXING



Passive tracers dynamics are identical to other viscous cases

Bodies move toward viscosity minima, this is consistent with work done by (Li, McKinley and Ardekani, 2015)

Viscosity gradient effects compete with geometry induced chaos



MOTION OF MULTIPLE ELLIPSOID IN INVISCID FLOW

 $\nabla^2 \phi = 0$

We know that the ϕ will have the form

$$\phi = \dot{q_1}\phi_1 + \dot{q_2}\phi_2 + \dots + \dot{q_n}\phi_n$$

The total energy of the fluid:

$$2\mathbf{T} = -\rho \iint \phi \frac{\partial \phi}{\partial \hat{n}} dS = \sum_{i,j=1}^{n} 2\mathbf{A}_{i,j} \dot{q}_i \dot{q}_j$$

$$\boldsymbol{A}_{i,j} = -\rho \iint \phi_i \frac{\partial \phi_j}{\partial \boldsymbol{\hat{n}}}$$





- Correlations in the chaotic motion of two bodies
- Diagonal lines represent times of correlations and perpendicular lines present times of anticorrelation between bodies
- Further Investigation underway







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CONCLUSION & FUTURE WORK

Conclusion

Our GISS solver is capable to capture very complex motion patterns of solidfluid.

Parametric mapping of the systems supports trend found by (Aref & Jones, 1993)

Vortex Shedding allows for the development of chaotic orbits

Future Work

Investigate the enhancement to mixing caused by the chaotic tumbling ellipsoids

Extent Kirchhoff Equation to include multiple ellipsoids

THANK YOU

For your attention