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# Single Bubble Collapse at Audible Frequencies and High Amplitudes

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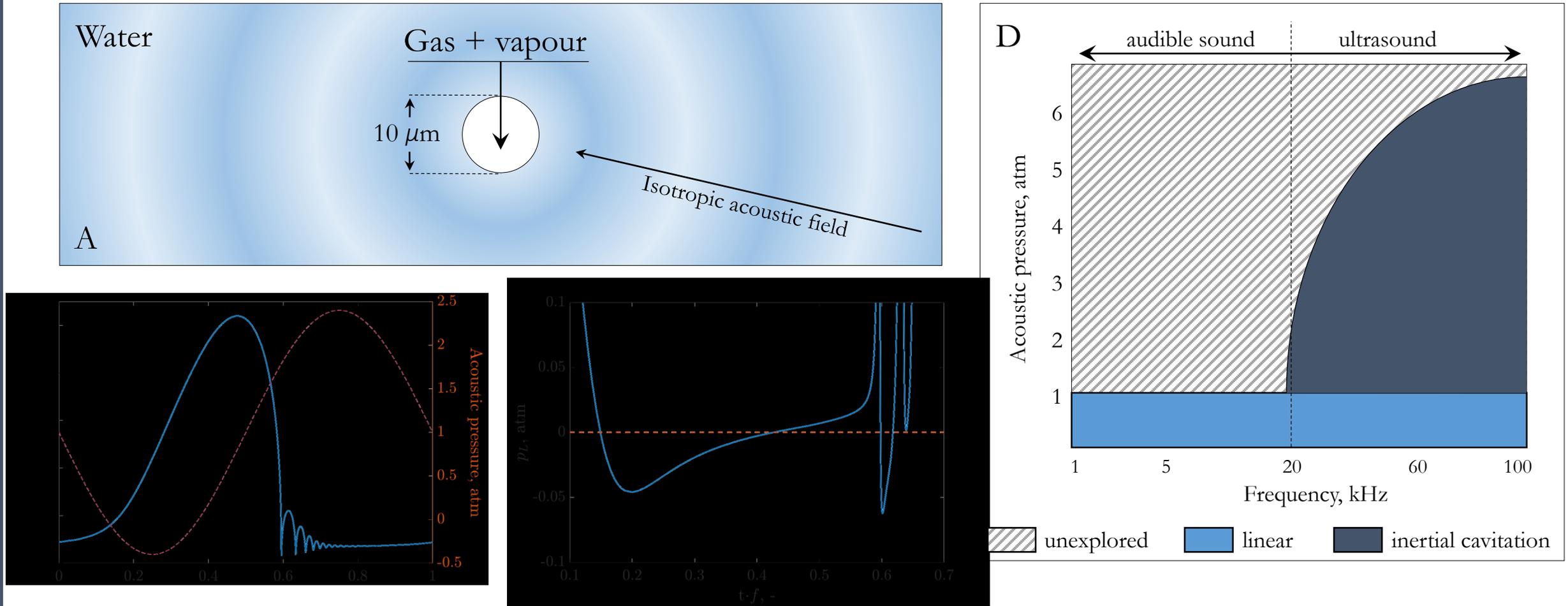
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# Phenomenon and motivation



**Figure 1.** A) Schematic of the system; B) Bubble dynamics; C) Liquid pressure; D) Unexplored parameters space.

# Governing equation

Bubble wall motion (Keller-Miksis equation with effect of bulk viscosity)

$$\left(1 - \frac{\dot{R}}{c_l}\right) R \ddot{R} + \frac{3}{2} \left(1 - \frac{\dot{R}}{3c_l}\right) \dot{R}^2 = \left(\frac{R}{\rho_l c_l} + \frac{\lambda_l + 2\mu_l}{\rho_l^2 c_l^2}\right) (\dot{p}_l - \dot{p}_\infty) + \left(1 + \frac{\dot{R}}{c_l}\right) (p_l - p_\infty)$$

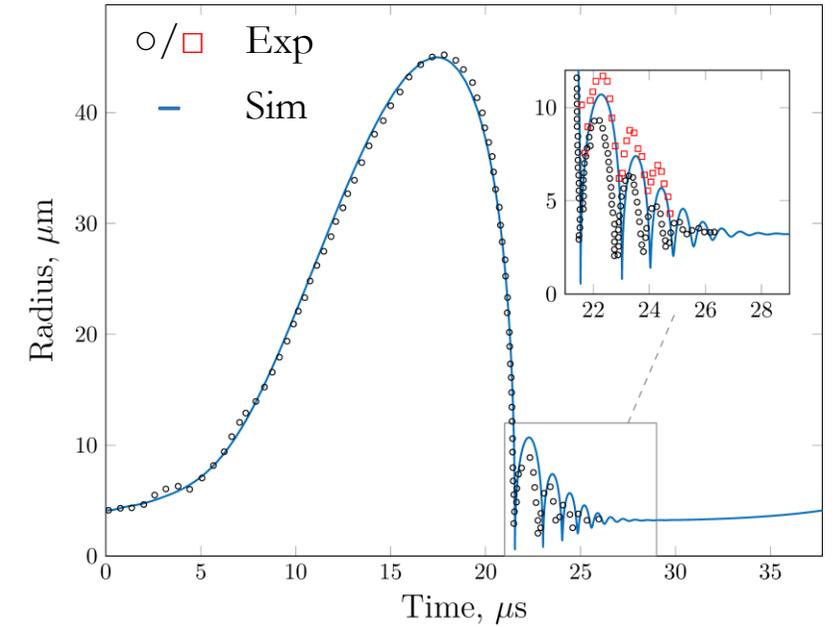
Gas temperature (quasi-equilibrium assumption)

$$\dot{T}_g = \frac{-T_g \frac{\partial p_g}{\partial T_g} dV + S \kappa_g \frac{\partial T_g}{\partial r} \Big|_{r=R} - S \sum_{i=1}^{N_s} \Psi_i \left[ u_i^0(T_g) - u_i^0(T_g^i) \right] - V \sum_{j=1}^{N_R} r_j^{\text{net}} \Delta^r u_j^0}{n_{TOT} \bar{c}_{vg}}$$

Non-equilibrium phase change (Hertz-Knudsen equation)

$$\dot{m} = \frac{\alpha}{\sqrt{2\pi\mathcal{R}}} \left( \frac{p_{\text{sat}}(T_{li})}{\sqrt{T_{li}}} - \frac{\Gamma x_{H_2O} p_g}{\sqrt{T_{gi}}} \right)$$

[1] R. Löfstedt, B. P. Barber, and S. J. Putterman, Physics of Fluids A: Fluid Dynamics 5, 2911 (1993).



**Figure 2.** Model validation against experimental data of bubble radius from Löfsted et al. [1]

## Physical properties

*Gas phase*

Redlich – Kwong EOS

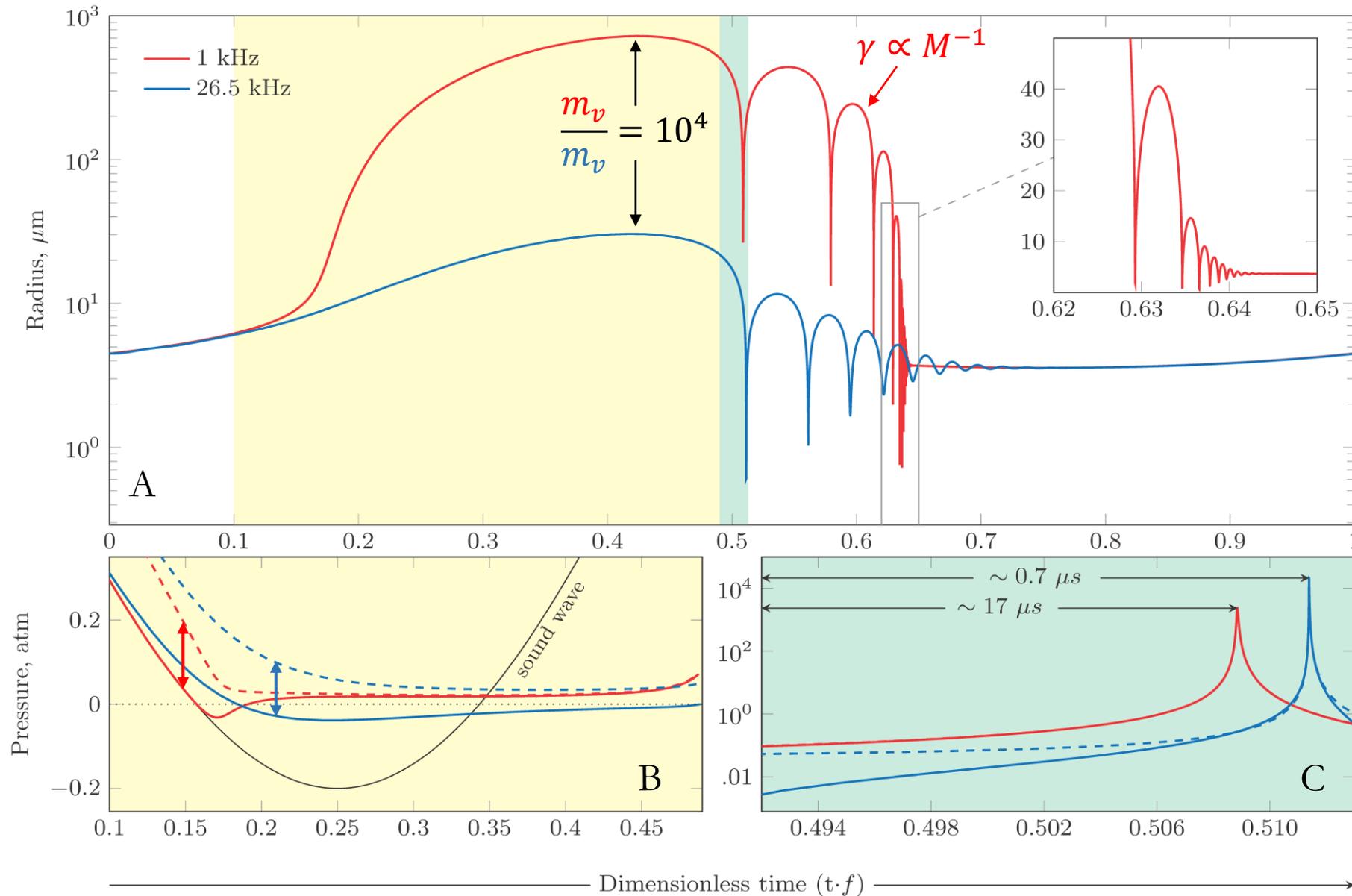
NASA polynomials

Chapman – Enskog theory

*Liquid phase*

IAPWS R6-95 EOS

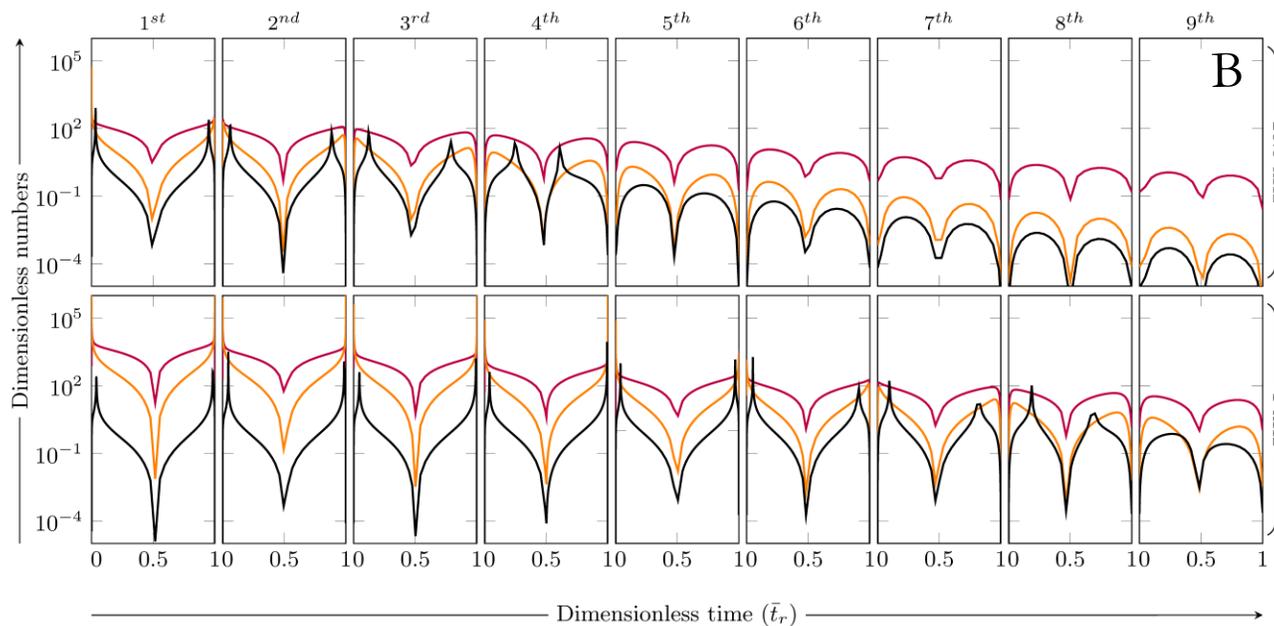
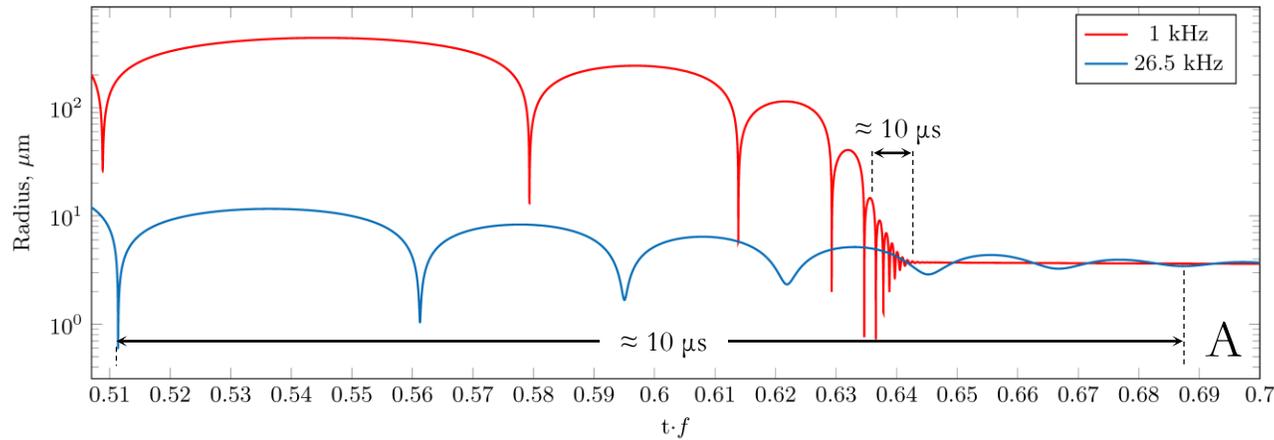
# Low vs. high frequency: the main collapse



**Figure 3.**

- A) Bubble radius dynamics at low and high frequency;
- B) Liquid and gas pressures during bubble expansion;
- C) Liquid and gas pressures at collapse.

# Low vs. high frequency: the rebounds



## Key aspects

1. Large bubble prevents viscosity and surface tension damping effect.
2. For same size, oscillations have same time scale regardless of driving frequency.

## Bubble wall dimensionless numbers

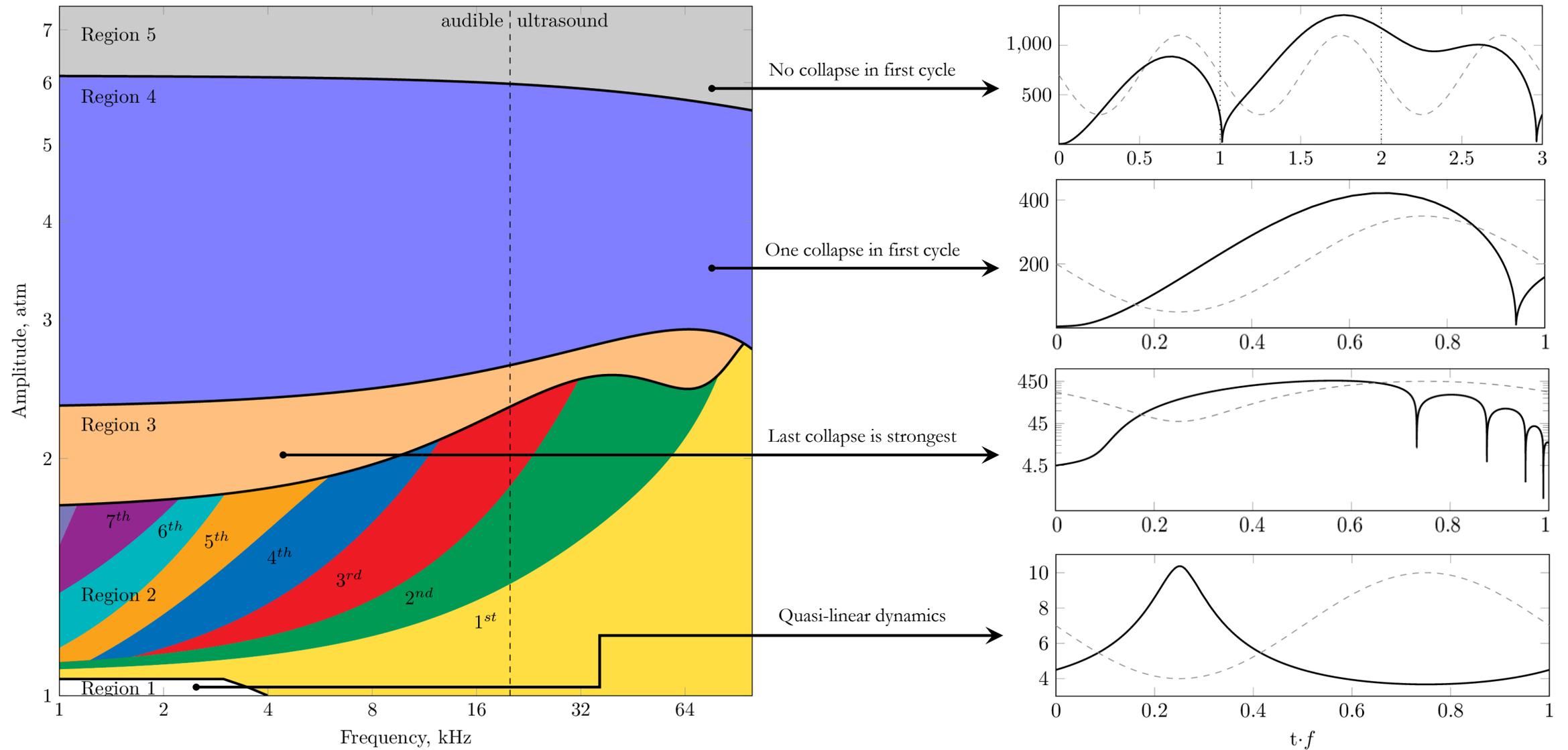
Reynolds:  $Re = \frac{R|\dot{R}|\rho_l}{\mu_l}$

Weber:  $We = \frac{R|\dot{R}|^2\rho_l}{\sigma_l}$

Pressure:  $\Pi_g = \frac{|\dot{R}|^2\rho_l}{p_g}$

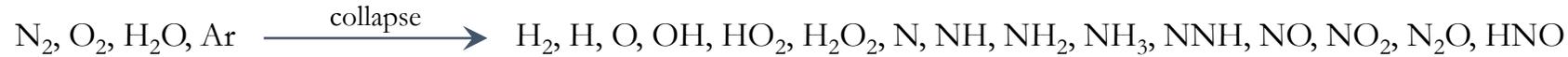
**Figure 4.** A) Inset of the rebounds following the main collapse. B) Dimensionless numbers evolution for each rebound.

# Different dynamics

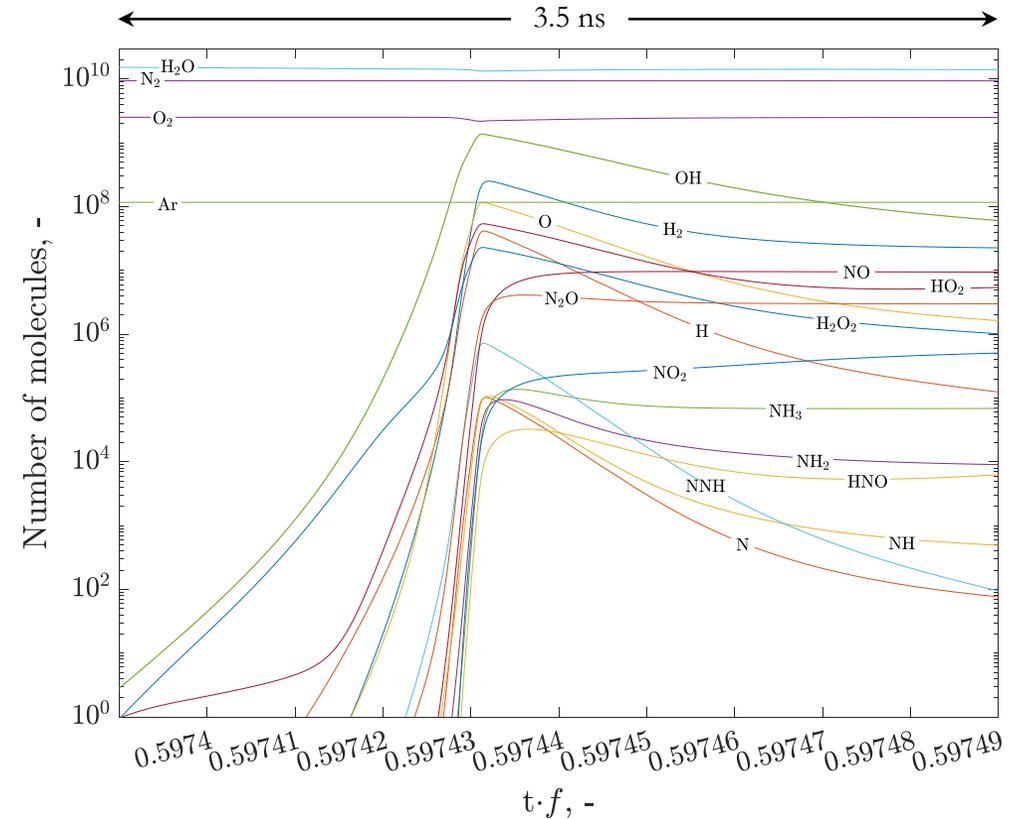
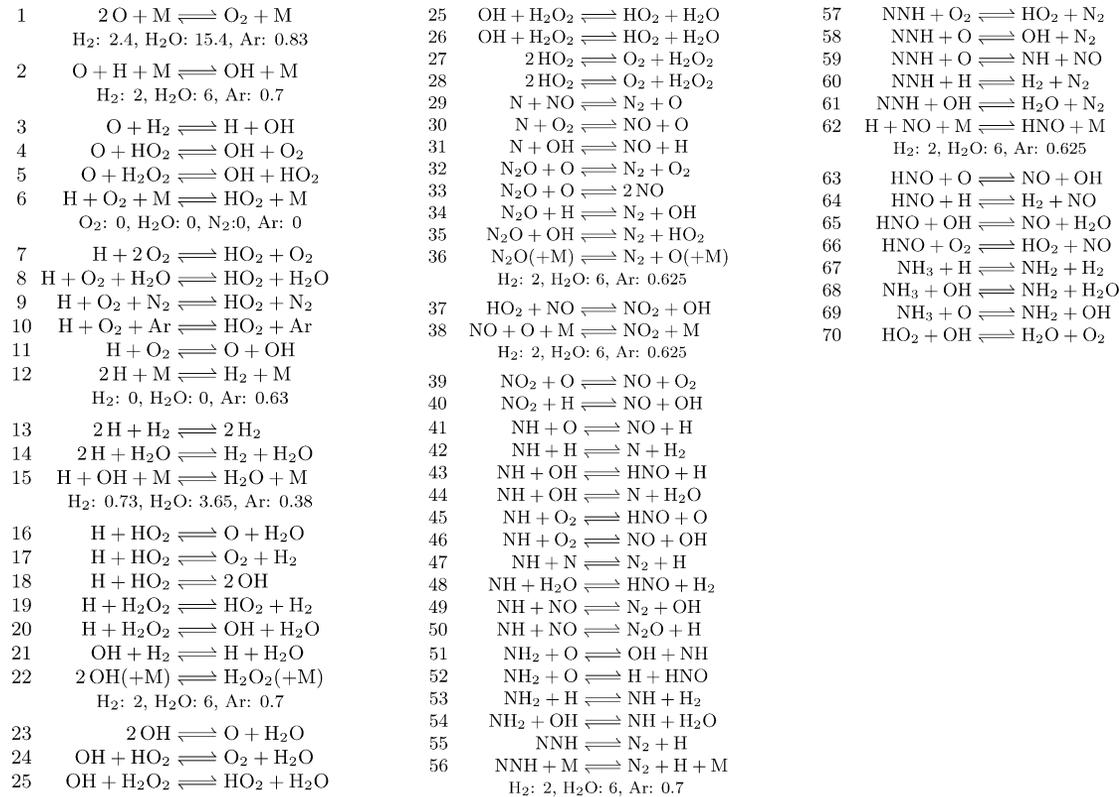


**Figure 5.** Division of the parametric space by typology of bubble dynamics within one acoustic cycle.

# Gas phase chemistry

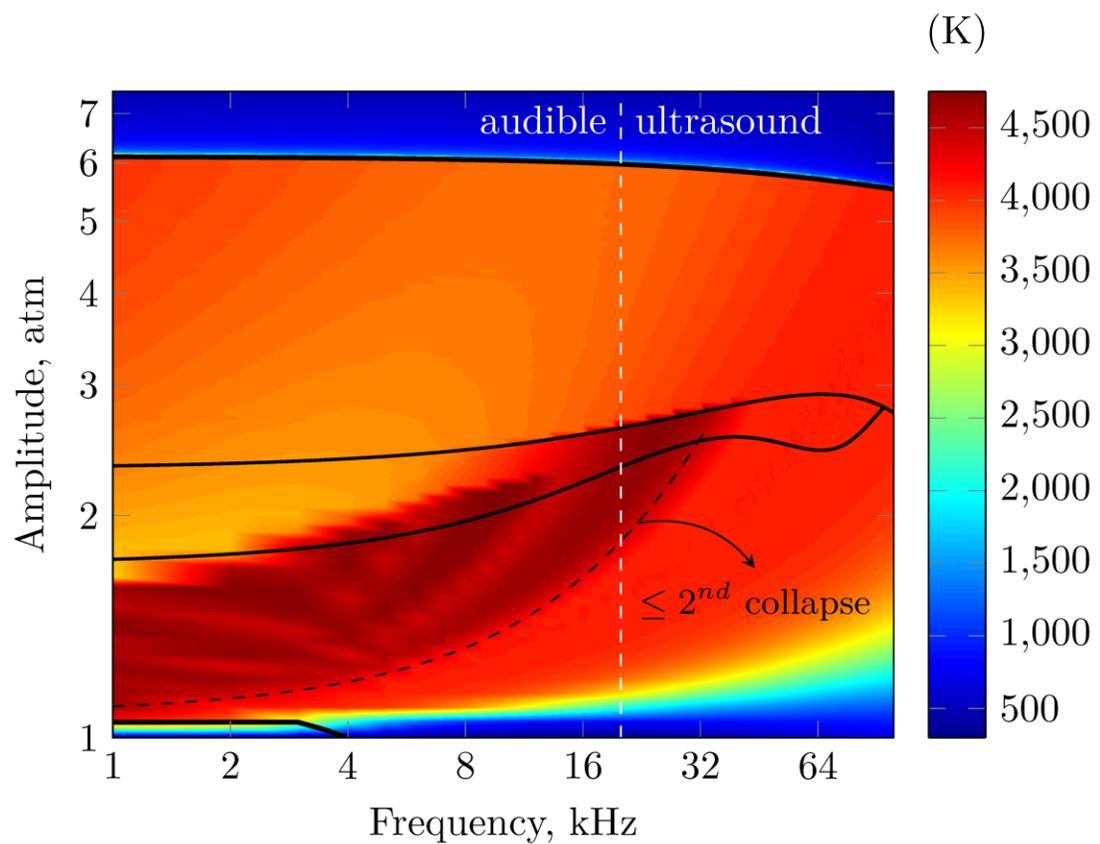


## GRI 3.0 combustion mechanism

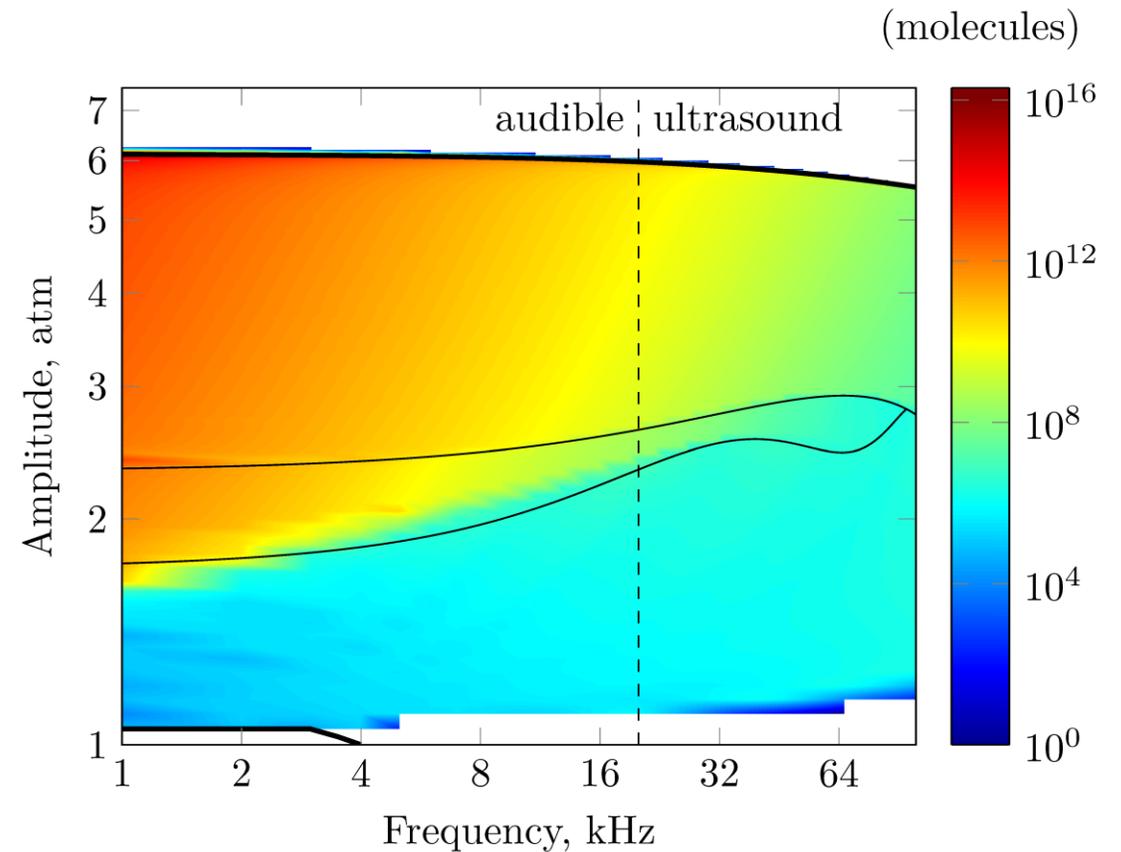


**Figure 6.** Simulated reaction network kinetics for a 5 μm air bubble collapsing at 26 kHz and 1.4 atm.

# Temperature and chemistry



A) Maximum temperature during one acoustic cycle

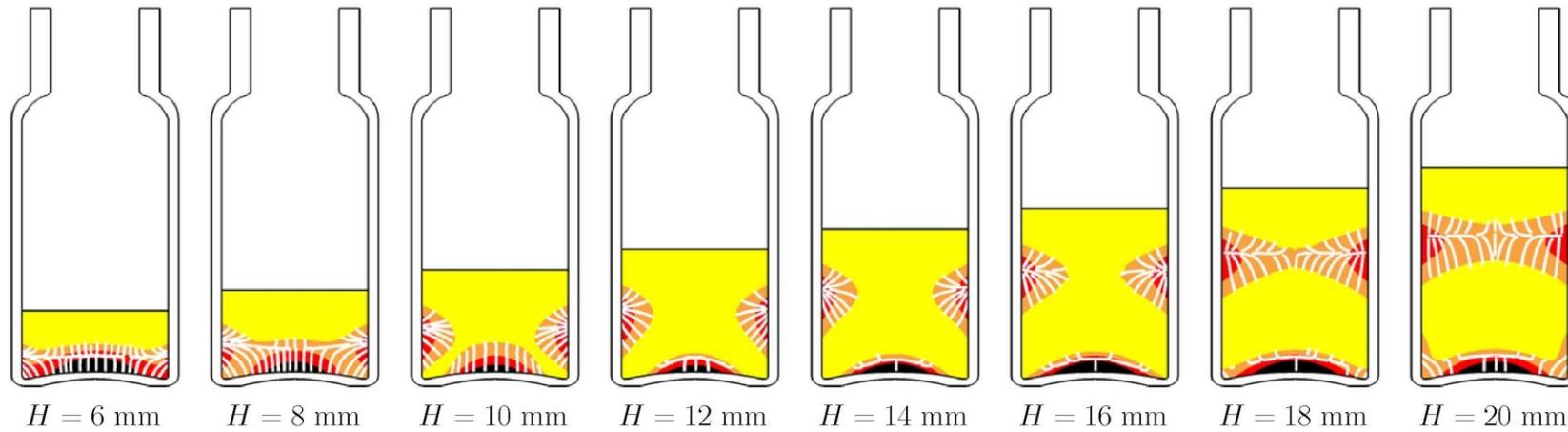


B) Number of OH radicals in the molecules at the end of the cycle

**Figure 7.** Maximum temperature and produced amount of hydroxyl radicals within one acoustic cycle in the explored parametric space.

# Can we use low frequency sound in sonochemistry?

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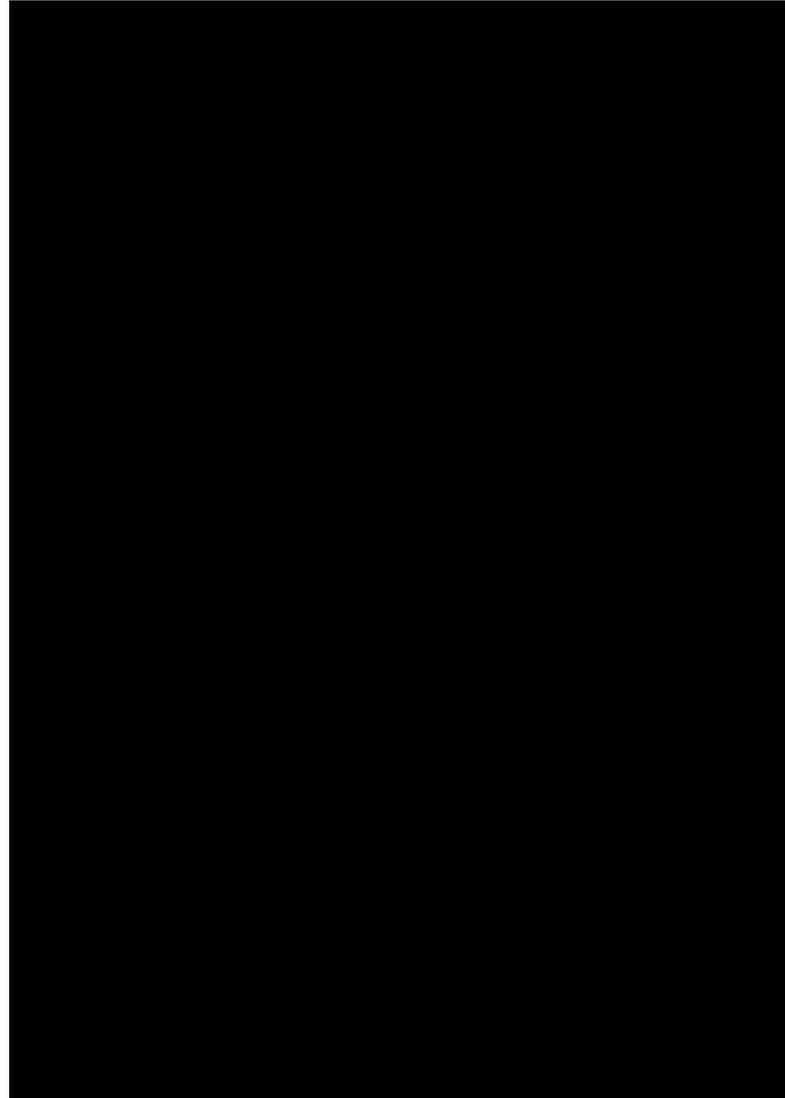
**Figure 8.** Change of acoustic field in bubbly water with size of propagating medium<sup>1</sup>.

- Implementation of rectified diffusion in the model equations.
- Evaluation of Rayleigh-Taylor and parametric instabilities.

<sup>1</sup>Louisnard, O., et al., 2015. Prediction of the acoustic and bubble fields in insonified freeze-drying vials. *Ultrasonics sonochemistry*, 26, pp.186-192.

# Future work

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**Video 1.**  
Appearance  
of acoustic  
streamers  
in a 18 kHz  
acoustic  
standing  
wave.